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# Well GeHP detector calibration for environmental measurements using reference materials

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## Abstract

A well-type detector installed in the Modane underground Laboratory (LSM) can combine both low background and high detection efficiency and it is well suited for the analysis of small amounts of environmental samples. Reference materials such as IAEA-447 (moss-soil), IAEA-RG-Th1 and IAEA-RG-U1 were used for the detector calibration, owing to a chemical composition close to those of the environmental samples. Nevertheless, the matrix effects and the true coincidence summing effects must be corrected from the full energy peak efficiency (FEPE). The FEPE was performed for a wide range of energy by a semi-empirical method using Monte Carlo simulation (MCNP6), intended for environmental measurements such as lake sediments dating. In the well geometry, the true coincidence summing effects could be very important and correction factors have been computed in three different ways.

**Keywords:** Well-type detector; Full energy peak efficiency; True coincidence summing; MCNP6; Efficiency correction

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## 1. Introduction

The detector full energy peak efficiency (FEPE) calibration is always required to reach accurate measurements and still represents a subject of considerable interest for

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4 the gamma spectrometry [1, 2, 4–10]. In this work, the measurements were carried  
 5 out with a well high-purity germanium (HPGe) detector, situated in the underground  
 6 Laboratory of Modane (LSM, located along the Frejus Tunnel Road in Savoy) where the  
 7 reachable sensitivity depends on the ultra-low background. The FEPE calibration of the  
 8 well detector [1, 2, 4, 7, 8, 10] is required for environmental measurements, mainly the  
 9 determination of fallout radionuclides in lake sediments such as  $^{210}\text{Pb}$  (half-life 22.3 years)  
 10 and  $^{137}\text{Cs}$  (half-life 30.05 years) for dating methods. To determine the FEPE function,  
 11 standard calibration sources, with the same size and composition as the samples, should  
 12 be used to have the same matrix effects (concerning mainly the auto-absorption due to  
 13 the density and the chemical composition) that will be corrected. Reference materials  
 14 such as IAEA-447 (moss-soil), RG-Th1 (thorium ore) and RG-U1 (uranium ore) were  
 15 used in the calibration process. Therefore, an important correction applied in close  
 16 measurements for this kind of detector, is related to the true coincidence summing effect  
 17 (TCS). TCS occurs when two (or more) emitted gamma or X-rays from a nucleus are  
 18 simultaneously detected within the resolving time of the gamma spectrometer system.  
 19 The magnitude of this effect depends on the detector efficiency (including the specific  
 20 source-detector geometry for a well detector) and the decay-scheme parameters. For  
 21 the concerned nuclei, TCS usually results in lower full-energy peak areas. In order to  
 22 compensate this loss of counts, a suitable correction must be performed. TCS correction  
 23 factors were computed for  $^{214}\text{Bi}$  (half-life 20 min), using secular equilibrium between  
 24  $^{226}\text{Ra}$  (half-life 1600 years) and their progenies, especially  $^{214}\text{Bi}$  and  $^{214}\text{Pb}$  with their  
 25 free lines of TCS effect at 295.22 keV and 351.93 keV. A simple way to get the correction  
 26 factor is to compare their activities, for example with the Genie 2000 software [11], which  
 27 should be equal.

28 A comprehensive study of the efficiency calculation and calibration verification of  
 29 the well spectrometer was performed by a semi-empirical method using Monte-Carlo  
 30 simulation, where the key element for calibration is the accurate knowledge of the physical  
 31 and geometrical characteristics of the whole detector, such as length and diameter of the  
 32 Ge crystal, thickness of the dead layer and more generally features of other components  
 33 (endcap, crystal holder, insulators, ...). A detector model was created for the efficiency  
 34 calculation, using the MCNP6 code [12]. This allowed to achieve a better accuracy for the

activity measurements of samples with unusual shapes, where experimental calibration with standard sources appears to be difficult.

## 2. Materials and methods

### 2.1. Well HPGe detector

The spectrometer used in this work is a Canberra High-Purity Germanium (HPGe) well detector, model GCW4021 under a serial number b07035, with a relative efficiency of 40%, an active volume of 238 cm<sup>3</sup> and a nominal FWHM of 1.27 keV at 122 keV (<sup>57</sup>Co) and 2.03 keV at 1.33 MeV (<sup>60</sup>Co). These features were supplied by the manufacturer. It works coupled to a DSA-1000 Canberra multichannel analyser and provides a maximum efficiency for small samples, because the sample is virtually surrounded by the active detector material.

The Canberra well detector is fabricated with a blind hole rather than a through hole, leaving at least 15 mm of active detector thickness at the bottom of the well. The counting geometry therefore approaches  $4\pi$  sr.

The well insert in the end-cap is made of low background (LB) aluminium with a side-wall thickness of 1.5 mm and a 1 mm thick bottom. The ion boron implanted contact on the detector element is negligibly thin compared to 0.5 mm of aluminium, so this kind of detector has intrinsically a good response at low energy, down to 20 keV [13]. The detector is shielded with a foil of electrolytic copper (3 mm thick) and lead (12 cm of low activity lead  $< 50$  Bq kg<sup>-1</sup> and 3 cm of very low activity lead  $< 10$  Bq kg<sup>-1</sup>).

A completed description of the equivalent detector model is represented in Fig. 1, where the dimensions are taken from the Canberra handbook.

Table 1 summarizes the values specified by Canberra for the geometric features of the well detector shown in Fig. 1.

### 2.2. Monte-Carlo simulations

The characterisation of the detector usually combines experimental measurements and Monte-Carlo simulations to calculate with accuracy the detector efficiency.

An initial model of the detector was performed by using the nominal dimensions and features provided by the manufacturer and then was implemented in the Monte-Carlo

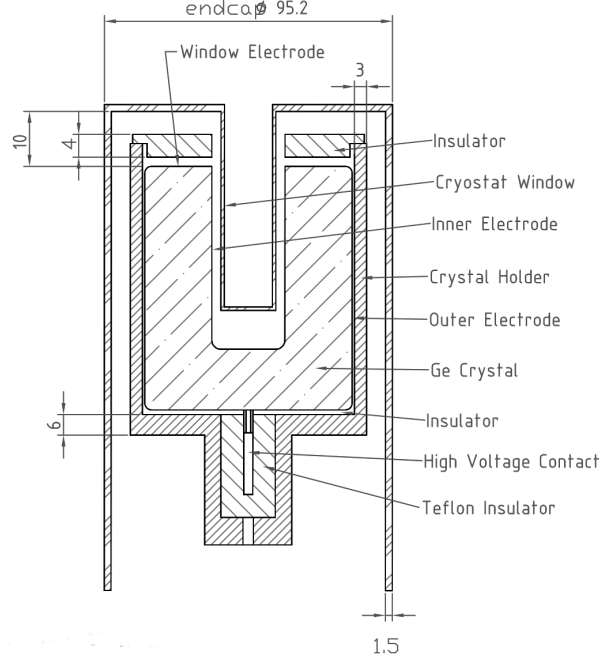


Figure 1: Well detector longitudinal section.

Parameters	Value (mm)
Outer electrode thickness	0.9
Inner electrode thickness	$0.3 \times 10^{-3}$
Window electrode thickness	0.9
Cristal diameter	68
Cristal length	68
Core hole diameter	17
Endcap hole	diameter 11 mm and depth 40 mm
Core hole depth	35
Cryostat window material	LB aluminium 1 mm thick
Endcap material	LB aluminium 1.5 mm thick
Crystal holder	LB copper

Table 1: Well detector parameters.

code MCNP6. This detector model must be checked by comparing the efficiency curve provided by MCNP6 with the experimental one obtained for source-detector geometry and reference materials in a wide energy range. The detector model should be approved if the calculated efficiencies are in a good agreement with the experimental values according to a level of acceptable uncertainty. Note that MCNP6 intrinsically does not take into account the TCS effect.

Fig. 2 shows the detector model obtained by MCNP6. The detector resolution was taking into account through the GEB (Gaussian Energy Broadening) card.

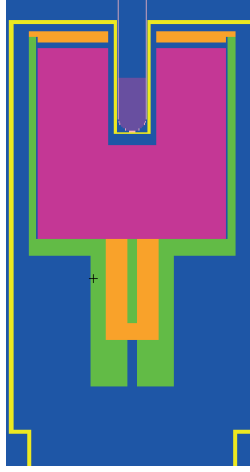


Figure 2: 2-D representation of the well detector MCNP 3-D model.

### 2.3. IAEA447 Standard

The experimental efficiency calibration was carried out in the 46.54–2614.51 keV energy range using the IAEA-447 standard (moss-soil). The milled material was sieved to obtain a maximum particle size distribution of 150  $\mu\text{m}$ . The material density was measured in 5 test portions and found to be  $1.03 \pm 0.05 \text{ g cm}^{-3}$  [14].

The certified values used to evaluate the activities of radionuclides were established on the basis of results reported by the IAEA Terrestrial Environment Laboratory in Seibersdorf, Austria [14].

Radionuclides from this sample and their activities estimated at the date of our experiment are represented in Table 2.

Radionuclides	Certified Values (Bq kg <sup>-1</sup> )	Uncertainty (Bq kg <sup>-1</sup> )
<sup>137</sup> Cs	383.46	10.00
<sup>210</sup> Pb	378.65	0.02 × 10 <sup>3</sup>
<sup>210</sup> Po	378.65	10.00
<sup>212</sup> Pb	8.28	1.50
<sup>226</sup> Ra	25,05	2.00
<sup>228</sup> Ac	15,49	2.00
<sup>234</sup> U	21.80	0.80
<sup>238</sup> U	22.20	0.80
<sup>238</sup> Pu	0.15	0.02
<sup>239–240</sup> Pu	5.30	0.20
<sup>40</sup> K	550	0.02 × 10 <sup>3</sup>
<sup>90</sup> Sr	4.50	0.30
<sup>232</sup> Th	37.30	20
<sup>241</sup> Am	2.20	0.20
<sup>241</sup> Pu	6.45	1.00

Table 2: Estimated values of activities by Darwin software from IAEA-447 certified ones.

#### 2.4. IAEA-RGU-1 and IAEA-RGTh-1 Standards

Both IAEA-RGU-1 and IAEA-RGTh-1 reference materials were prepared, on behalf of the International Atomic Energy Agency by the Canada Centre for Mineral and Energy Technology, by dilution of respectively a uranium ore BL-5 (7.09% U) and a thorium ore OKA-2 (2.89% Th, 219 µg U/g) with floated silica powder of similar grain size distribution. BL-5 has been certified for uranium, <sup>226</sup>Ra and <sup>210</sup>Pb, confirming that it is in radioactive equilibrium. The agreement between radiometric and chemical measurements of thorium and uranium in OKA-2 shows that both series are in radioactive equilibrium [15].

The activities of radionuclides from IAEA-RG-Th1 and IAEA-RG-U1 are shown in Table 3.

IAEA-RG-Th1		IAEA-RG-U1	
Radionuclide	Activity (Bq kg <sup>-1</sup> )	Radionuclide	Activity (Bq kg <sup>-1</sup> )
<sup>232</sup> Th	3250	<sup>232</sup> Th	< 4
<sup>235</sup> U	3.6	<sup>235</sup> U	238
<sup>238</sup> U	78	<sup>238</sup> U	4940
<sup>40</sup> K	6.3	<sup>40</sup> K	< 0.68

Table 3: IAEA-RG-Th1 (thorium ore) and IAEA-RG-U1 (uranium ore) certified values of activities.

### 2.5. Experimental

All measurements in this work were performed at the LSM. The laboratory is shielded from cosmic radiation by 1700 m of rocks, equivalent to 4400 m of water, and an air flushing without radon (generated from the radon trapping facility at the LSM) is done into the measurement cell, inside the lead shielding. Thus the background rate between 20 keV and 2 MeV is of 23 counts h<sup>-1</sup>. Three reference materials cited above were used in containers with the following features:

- IAEA-447 in a PE tube of 28 mm of height and 1.341 g of weight;
- IAEA-RG-Th1 in a PE tube of 28 mm and 1.4 g;
- IAEA-RG-U1 in a PE tube of 28 mm and 1.849 g.

The reference material samples are well sealed to ensure its air-tightness, so secular equilibrium between <sup>226</sup>Ra and <sup>210</sup>Pb can be reached after 20 days with 95%. In environmental samples measurement, the activity levels are low, so if statistically significant results require long count periods, they are necessary until several days in this work.

## 3. Results and discussion

### 3.1. FEP efficiency

For the analysis of the IAEA-447 standard, a correction factor is generally required for the spectral interfering  $\gamma$ -rays to determine the net areas of the analytical peaks, because some of those interferences often might contribute to the analytical peaks of



interest [16]. The nuclide identification was performed using our library containing radio-isotopes presented in the IAEA-447 sample where  $^{238}\text{U}$ ,  $^{232}\text{Th}$  progenies and  $^{40}\text{K}$  were identified. About  $^{234}\text{Th}$  gamma emission at 92.38 keV, there is a single peak in the spectral region 92-93 keV resulted from two energies (92.56 keV and 92.78 keV), where the total emission probability was taken into account (the net total peak area of this unresolvable multiplet is accounted for without deconvolution).

Table 4 lists the main correction factors for some radionuclides peaks [16].

Radionuclide	Nuclides in the peak	Energy (keV)	Proportion in the peak (%)
$^{234}\text{Th}$	$^{234}\text{Th}$	63.28	98.2
	$^{232}\text{Th}$	63.81	1.8
$^{226}\text{Ra}$	$^{226}\text{Ra}$	186.21	57.1
	$^{235}\text{U}$	185.72	42.8
$^{212}\text{Pb}$	$^{212}\text{Pb}$	238.63	62.4
	$^{214}\text{Pb}$	242.00	31.7
	$^{224}\text{Ra}$	240.99	5.9
$^{40}\text{K}$	$^{40}\text{K}$	1460.82	94.8
	$^{228}\text{Ac}$	1459.14	5.2

Table 4: Interference factors of such radionuclides in IAEA-447 [16].

Table 5 compares the experimental and simulated efficiencies for the IAEA-447 standard for the most intense peak of radionuclides. To calculate the efficiency, we have used the following equation:

$$\varepsilon = \frac{N_{net}}{A(Bq) \times I_{\gamma} \times t(s)} \quad (1)$$

where  $A$  is the initial activity carried out with the Darwin software [17] taking into account the radioactive affiliations,  $N_{net}$  is the number of counts under the net peak area,  $I_{\gamma}$  is the probability of gamma emission and  $t$  the acquisition time.

A number of analytical functions describing the dependence of the FEPE as a function of the energy have been proposed by several authors [18, 19]. The efficiency function used

Radionuclide	Energy (keV)	Intensity (%)	$\varepsilon_{\text{exp}}$	$\varepsilon_{\text{sim}}$	Ratio ( $\varepsilon_{\text{exp}}/\varepsilon_{\text{sim}}$ )
$^{210}\text{Pb}$	46.54	4.25	$0.422 \pm 0.006$	$0.457 \pm 0.001$	$0.92 \pm 0.01$
$^{241}\text{Am}$	59.54	35.90	$0.482 \pm 0.060$	$0.536 \pm 0.044$	$0.90 \pm 0.23$
$^{234}\text{Th}$	63.28	4.80	$0.492 \pm 0.040$	$0.539 \pm 0.068$	$0.91 \pm 0.22$
	92.37	2.81	$0.531 \pm 0.02$	$0.561 \pm 0.067$	$0.95 \pm 0.16$
$^{226}\text{Ra}$	186.21	3.55	$0.450 \pm 0.02$	$0.474 \pm 0.038$	$0.95 \pm 0.12$
$^{212}\text{Pb}$	238.63	43.60	$0.375 \pm 0.005$	$0.384 \pm 0.069$	$0.98 \pm 0.19$
$^{214}\text{Pb}$	295.22	18.5	$0.294 \pm 0.010$	$0.311 \pm 0.025$	$0.95 \pm 0.11$
	351.93	35.6	$0.253 \pm 0.008$	$0.263 \pm 0.021$	$0.96 \pm 0.11$
$^{208}\text{Tl}$	583.19	85.00	$0.144 \pm 0.007$	$0.149 \pm 0.025$	$0.97 \pm 0.22$
$^{214}\text{Bi}$	609.31	45.49	$0.077 \pm 0.007$	$0.142 \pm 0.011$	$0.54 \pm 0.17$
	1120.29	14.90	$0.039 \pm 0.010$	$0.071 \pm 0.005$	$0.55 \pm 0.32$
$^{137}\text{Cs}$	661.66	84.99	$0.125 \pm 0.001$	$0.130 \pm 0.004$	$0.96 \pm 0.04$
$^{228}\text{Ac}$	911.20	25.80	$0.095 \pm 0.004$	$0.093 \pm 0.012$	$1.02 \pm 0.17$
$^{40}\text{K}$	1460.82	10.60	$0.053 \pm 0.001$	$0.052 \pm 0.002$	$1.02 \pm 0.18$
$^{208}\text{Tl}$	2614.51	99.75	$0.026 \pm 0.001$	$0.027 \pm 0.004$	$0.96 \pm 0.05$

Table 5: Experimental and calculated efficiencies for the IAEA-447 standard.

in this work has the form of logarithmic positive power transferred series, which has been proposed in [20, 21]:

$$\ln \varepsilon = a_0 + a_1 \ln E + a_2 \ln^2 E + a_3 \ln^3 E + a_4 \ln^4 E \quad (2)$$

Fig. 3 shows the experimental and simulated efficiencies as a function of the energy, where both were fitted using a fourth order polynomial from Eq. 2, because it is the adequate order which groups the best statistical parameters, such as trust and residues factor which must be the smallest in the sense of least squares method.

From Table 5, we can see that the simulated values are always greater than experimental values, excepted for  $^{228}\text{Ac}$  at 911.2 keV and for the primordial isotope of potassium  $^{40}\text{K}$  at 1460.82 keV, due to the presence of this isotope everywhere and chiefly from the human radioactivity. A good agreement is found to be within 10% between almost

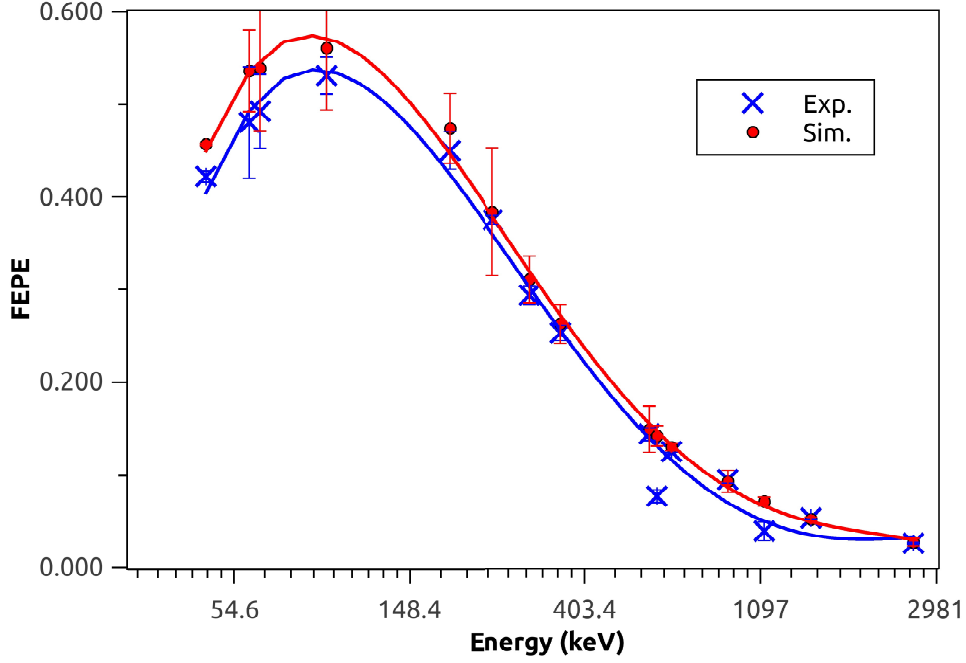


Figure 3: Comparison between experimental and calculated values of the FEP efficiency for the IAEA-447 standard.

all experimental and simulated values as shown in Table 5 and Fig. 3: experimental and simulated efficiencies are very close, starting from 92.37 keV line for  $^{234}\text{Th}$  with the ratio  $(0.95 \pm 0.16)$ . In the low energy range where the self-attenuation phenomenon is not insignificant, a correction in this case should be applied. For some lines with a low number of counts, the uncertainties are important, due to the behaviour of the analysis software towards weak peaks. The most important differences are for  $^{214}\text{Bi}$  at the two main lines 609.31 keV and 1120.29 keV, where only 54-55% of counts were detected. In this case, the TCS effect decreases the count number by summing those energies with other in succession.

### 3.2. True Coincidence Summing correction

The true coincidence summing effect was observed for radio-isotopes which have a complex decay scheme, such as  $^{214}\text{Bi}$  which can be used in the  $^{226}\text{Ra}$  activity evaluation, required in the  $^{210}\text{Pb}$  dating method of lake sediments. For  $^{214}\text{Bi}$ , we have considered

height major lines (with the most important intensities) from its numerous  $\gamma$  radiations. We assumed that  $^{226}\text{Ra}$  is in secular equilibrium with its progenies, which implies that they have the same activity, and we have compared the  $^{214}\text{Bi}$  activity with that of  $^{214}\text{Pb}$ , which is free from TCS at 351.9 keV and slightly affected with 1.9% at 295.38 keV [8].

Table 6 shows the TCS correction factors calculated for the major lines of  $^{214}\text{Bi}$  in three different ways:

- Activity correction: it was given by the ratio between the raw activity of every line of  $^{214}\text{Bi}$  and the activity of  $^{214}\text{Pb}$ , used as a reference and equal to  $23.85 \pm 0.02$  Bq  $\text{kg}^{-1}$ ;
- TCS factor (Genie 2000): it was determined from the FEPE experimental fitting curve given by Genie 2000;
- Ratio  $\varepsilon_{\text{exp}}/\varepsilon_{\text{sim}}$  : that is the ratio between the experimental efficiency and the efficiency calculated from the MCNP6 simulations.

Energy (keV)	Intensity (%)	Activity (Bq.kg $^{-1}$ )	Correction factors		
			Activity Correction	TCS factor (Genie 2000)	Ratio ( $\varepsilon_{\text{exp}}/\varepsilon_{\text{sim}}$ )
609.31	45.49	$12.72 \pm 0.02$	0.53	0.58	0.54
1120.29	14.91	$8.91 \pm 0.06$	0.37	0.57	0.55
1238.1	5.83	$13.05 \pm 0.11$	0.55	0.56	0.55
1377.7	3.97	$32.07 \pm 0.18$	1.34	1.40	1.27
1729.5	2.84	$75.07 \pm 0.27$	3.14	3.07	2.78
1764.5	15.31	$40.70 \pm 0.09$	1.70	1.75	1.61
1847.42	2.11	$69.94 \pm 0.29$	2.93	2.15	1.76
2204.1	4.91	$53.99 \pm 0.23$	2.26	2.38	2.20

Table 6: TCS correction factors for the  $^{214}\text{Bi}$  major lines.

We can observe an important under estimation at 609.31 keV and 1120.29 keV, because these lines are the most involved in the TCS with the other photons emitted by  $^{214}\text{Bi}$ , leading to an important loss of counts for these peaks. The difference between the

167 correction factor obtained at 1120.29 keV and the TCS factor and the ratio  $\varepsilon_{\text{exp}}/\varepsilon_{\text{sim}}$   
 168 is due to the weak count rate in this line. The correction factors determined in three  
 169 different ways are in good agreement and are close to those obtained in the literature  
 170 for the most of values [8, 22]. For both lines at 1764.5 keV and 2204.1 keV, there are  
 171 significant differences compared with the values from the literature.

172 Fig. 4 shows the values of efficiency after correction from TCS for  $^{214}\text{Bi}$  at 609.31  
 keV and 1120.29 keV.

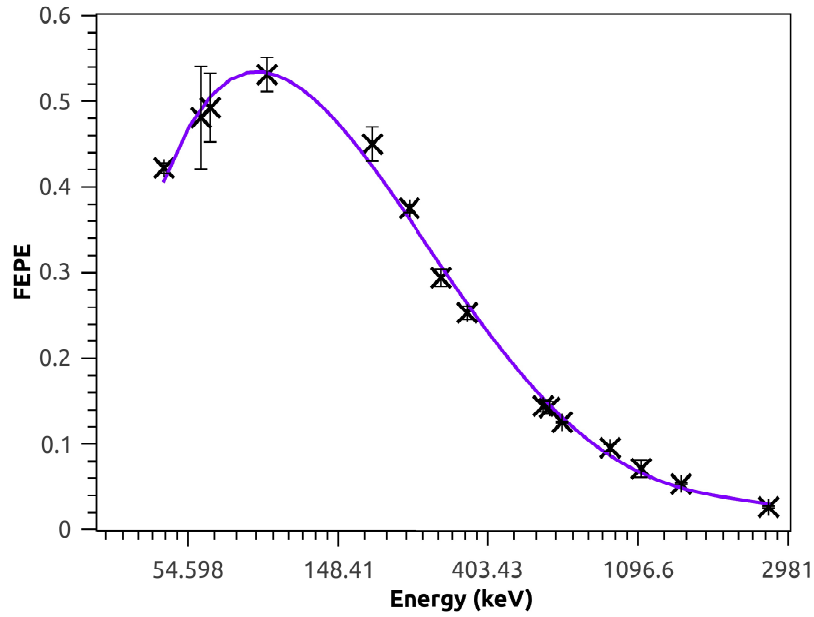


Figure 4: Experimental values of the FEPE for the IAEA-447 corrected from TCS.

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### 174 3.3. Density effect on TCS

175 We have compared the results of efficiencies obtained by IAEA-447, IAEA-RG-U1  
 176 and IAEA-RG-Th1 analysis with the common main lines of progenies from  $^{238}\text{U}$  and  
 177  $^{232}\text{Th}$ . The results are listed in Table 7, where efficiency ratios between two standards  
 178 are shown in the last column with their uncertainties for the same energy lines for the  
 179 three reference materials.

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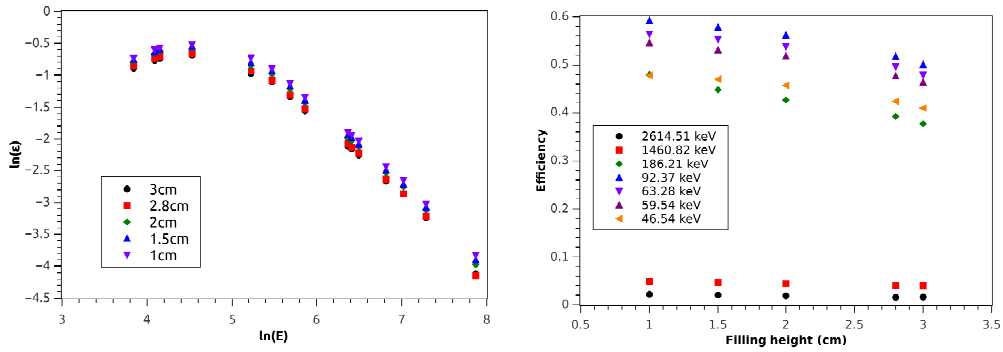
Radionuclide	Energy (keV)	$\varepsilon_{\text{IAEA-447}}$	$\varepsilon_{\text{RG-U1}}$	$\varepsilon_{\text{RG-Th1}}$	Ratio between efficiencies
$^{210}\text{Pb}$	46.54	$0.422 \pm 0.006$	$0.403 \pm 0.006$	—	$0.95 \pm 0.01$
$^{234}\text{Th}$	63.28	$0.492 \pm 0.040$	$0.485 \pm 0.007$	—	$0.99 \pm 0.10$
	92.37	$0.531 \pm 0.020$	$0.518 \pm 0.008$	—	$0.98 \pm 0.05$
$^{226}\text{Ra}$	186.21	$0.450 \pm 0.020$	$0.425 \pm 0.006$	—	$0.94 \pm 0.06$
$^{212}\text{Pb}$	238.63	$0.375 \pm 0.005$	—	$0.369 \pm 0.020$	$0.98 \pm 0.07$
$^{214}\text{Pb}$	295.22	$0.294 \pm 0.010$	$0.290 \pm 0.004$	$0.98 \pm 0.05$	
	351.93	$0.253 \pm 0.008$	$0.241 \pm 0.003$	—	$0.95 \pm 0.04$
$^{208}\text{Tl}$	583.19	$0.144 \pm 0.007$	—	$0.103 \pm 0.006$	$0.72 \pm 0.10$
$^{214}\text{Bi}$	609.31	$0.077 \pm 0.007$	$0.062 \pm 0.001$	—	$0.80 \pm 0.10$
	1120.29	$0.039 \pm 0.01$	$0.03 \pm 0.001$	—	$0.77 \pm 0.27$
$^{228}\text{Ac}$	911.2	$0.095 \pm 0.004$	—	$0.082 \pm 0.005$	$0.86 \pm 0.10$
$^{208}\text{Tl}$	2614.51	$0.026 \pm 0.001$	—	$0.014 \pm 0.001$	$0.54 \pm 0.10$

Table 7: Comparison between IAEA-447, IAEA-RG-U1 and IAEA-RG-Th1 results.

FEPE values of radionuclides presented in the IAEA-447 reference material are always greater than those obtained by IAEA-RG-Th1 and IAEA-RG-U1, because the TCS effect increases with activity where IAEA-RG radionuclides have greater activities than those of IAEA-447. The results show a slight discrepancy between common lines of IAEA reference materials less than 10% in the range of low energies. This is due to the density difference between IAEA samples, where heavy elements such as thorium and uranium are presented in IAEA-RG, leading to a higher density than IAEA-447. This density impact is presented for all energies. The difference between efficiencies is low for the intermediate energy range. For higher energies, mainly those of  $^{214}\text{Bi}$ ,  $^{208}\text{Tl}$ , and  $^{228}\text{Ac}$ , differences are greater and vary between 14% for 911.2 keV of  $^{228}\text{Ac}$  and reach 46% for  $^{208}\text{Tl}$  at 2614.51 keV. For such environmental matrices which contain high Z elements (e.g. thorium, uranium, lead, ...) in significant quantities, the threshold at which matrix effects take place is moved to higher energies [23].

### 3.4. Influence of sample height on the FEP efficiency

We have further used MCNP6 to compute the FEPE for the IAEA-447 standard as a function of the sample height. The sample heights are 1, 1.5, 2, 2.8 and 3 cm and the results obtained are shown in Fig 5.a. The impact of the sample height on the efficiency is shown in Fig 5.b for the lowest energies (corresponding to  $^{210}\text{Pb}$ ,  $^{241}\text{Am}$  and  $^{234}\text{Th}$ ) and for the highest energies (corresponding to  $^{40}\text{K}$  and  $^{208}\text{Tl}$  at 2614.51 keV).



(a) Simulated values of FEP efficiency for several IAEA-447 standard heights. (b) The impact of the IAEA-447 standard heights on FEP efficiency.

Figure 5: Simulated values of FEP efficiency as function of the IAEA-447 standard heights.

The results show that the impact of the sample filling height is very important in the low energy range compared to the high energy range. The efficiency decreases with increasing in the sample filling height. Therefore for our future measurements, we will take into account the height of the sample, through a specific FEPE depending on the filling height.

## 4. Conclusion

The FEPE calibration of an HPGe well-type detector was performed in this work motivated by applications in environmental measurements, especially to determine the activities of radioactive fallout in the lake sediments such as  $^{210}\text{Pb}$ ,  $^{241}\text{Am}$  and  $^{137}\text{Cs}$ , used in dating methods. We have used the IAEA-447 (moss-soil) reference material as a standard to approach at best the matrix effects due to the chemical composition of the samples. A well detector model was implemented in MCNP6 with the features

provided by the manufacturer. The simulated values are greater than the experimental ones for most of the energy lines. The results show mostly a good agreement with respect to the experimental values and discrepancies are within 10%. This allows us to determine the efficiency calibration curve without an experimental work, and can be considered as an efficiency transfer model, that can be used in other investigations such as self-attenuation in samples. We have also calculated the true coincidence summing correction factors for the lines emitted in the complex decay chains of  $^{214}\text{Bi}$  with three manners, which have showed a good agreement with the literature values. In order to show matrix effects such as the density influence, we have compared the results from the FEPE obtained by three reference materials IAEA-447, IAEA-RG-Th1 and IAEA-RG-U1. The impact of the sample filling height was investigated. It was remarkable for low energies, where the efficiency decreases with increasing in filling heights. A reduction of the filling height would diminish the matrix effects [5]. Finally, the efficiency curve can be used in activities evaluation for the environmental measurements, but to improve accuracy, other corrections must be done such as self-attenuation, density and sample height filling effects.

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